Communication

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Vapor-Assisted Ex-Situ Doping of Carbon Nanotube towards Efficient and Stable Perovskite Solar Cells

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KEYWORDS : perovskite, solar cell, carbon nanotube, doping, stability
ABSTRACT

Single-walled carbon nanotubes (CNTs) has been considered as a promising material for a top electrode of perovskite solar cells owing to its hydrophobic nature, earth-abundance, and mechanical robustness. However, its poor conductivity, a shallow work function, and non-reflective nature have limited further enhancement in power conversion efficiency (PCE) of top CNT electrode-based perovskite solar cells. Here, we introduced a simple and scalable method to address these issues by utilizing an ex-situ vapor-assisted doping method. Trifluoromethanesulfonic acid (TFMS) vapor doping of the free-standing CNT sheet enabled tuning of conductivity and work function of the CNT electrode without damaging underneath layers. The sheet resistance of the CNT sheet was decreased by 21.3% with an increase in work function from 4.75 eV to 4.96 eV upon doping of TFMS. In addition, recently developed 2D perovskite-protected Cs-containing formamidium lead iodide (FACsPbI$_3$) technology was employed to maximize the absorption. Because of the lowered resistance, better energy alignment, and improved absorption, the CNT electrode-based PSCs produced a PCE of 17.6% with a $J_{SC}$ of 24.21 mA/cm$^2$, $V_{OC}$ of 1.005 V, and FF of 0.72. Furthermore, the resulting TFMS-doped CNT-PSCs demonstrated higher thermal and operational stability than bare CNT and metal electrode-based devices.

Keywords: perovskite, carbon nanotube, doping, stability, electrode
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spiro-MeOTAD carbon nanotube

TFMS acid vapor

ACS Paragon Plus Environment
INTRODUCTION

Perovskite solar cells (PSCs), reported by Miyasaka and colleagues for the first time in 2009, have progressed remarkably through a heated power conversion efficiency (PCE) race after development of their solid-state version. With a surge of research efforts, the certified PCE of the lab scale PSCs increased from 14.1% to 23.3% for last 5 years, which is on par with silicon and CIGS solar cells. Now the research efforts are being focused on enhancing the durability of the devices and reducing the production cost.

PSCs are usually fabricated on transparent conducting oxide (TCO) glass substrates where the perovskite layer is sandwiched between additional selective charge transporting layers. Typically, a metal electrode is thermally evaporated on top to serve as a counter electrode. These top metal electrodes, however, are known to substantially increase the process and material cost as expensive gold or silver layers are deposited under high vacuum. Furthermore, the metal electrodes are found to be not robust enough to render long-term stability of the PSCs as the metal ions migrate into the bulk of the device under an operational condition to react with the active materials, thus degrading the device performance.

As an alternative to the metal electrodes, researchers have incorporated single-walled carbon nanotubes (CNTs) into the PSC system. The CNT is a favorable electrode material owing to its hydrophobic nature, earth-abundance, and mechanical robustness. The application of CNT realizes CNT-based PSCs with good stability and versatility. Among the reported CNT-used PSCs devices, aerosol-synthesized CNT top electrode (or back electrode) replacing metals in PSCs has shown the most promising potential. The application of CNT as the top electrode substantially enhances the stability of PSCs by removing the ion migration, and drastically reduces the fabrication cost as it can be easily deposited onto devices by a simple mechanical
Despite such advantages, there are three factors limiting the PCE of the CNT top electrode-based PSCs that need to be addressed: 1) the work function of the CNT does not align with the perovskite, leading to loss in voltage potential 2) CNT top electrode is less conductive than the metal counterparts, limiting the fill factor (FF) of the devices 3) the CNT back electrode is not reflective that the devices cannot maximize the light harvesting short-circuit current.

In this work, we addressed those three issues by engineering the CNT electrode and photoactive layer. We tuned the work function of CNT electrodes and increased their conductivity by using a vapor-assisted doping of trifluoromethanesulfonic acid (TFMS). Furthermore, by incorporation of low bandgap 2D perovskite-protected Cs-containing formamidium lead iodide (2D/3D FACsPbI$_3$, $E_g$=1.48 eV), we enabled the harvesting of long-wavelength light, enhancing the photocurrent of the devices. By combining those technologies all into one device system, CNT electrode-based PSCs produced a PCE of 17.6% with a $J_{SC}$ of 24.21 mA/cm$^2$, $V_{OC}$ of 1.005 V, and FF of 0.72. The obtained PCE is the highest among the values reported from CNT top electrode-based PSCs. Moreover, the resulting CNT-PSCs exhibited higher high-temperature operational stability than those of the devices based on metal electrodes as well as conventional CNT electrodes.

RESULTS AND DISCUSSION

Chemical doping of carbon allotropes based electrodes have been widely adopted for enhancing the conductivity of the electrode and thus device performance. However, doping carbon top electrode in solar cells is an challenging task and there has been only one report to the best our knowledge. This is because acid solutions used for the doping is subject to damage the photoactive layers underneath. Here, we achieved a successful doping of CNT top electrode
without damaging the device by incorporating vapor-assisted *ex-situ* doping of highly effective and durable TFMS.\textsuperscript{25} The TFMS chemical \textit{p}-dopant is known to dope carbon electrodes more effectively with a much longer durability than conventional HNO\textsubscript{3}.\textsuperscript{20} By simply placing a sheet of CNT film on TFMS vapor for several tens of seconds, we can achieve the doping (Figure 1a). The degree of doping can be easily controlled by varying the exposure time of the CNT sheet to TFMS vapor (the doping time was varied from 10 to 50 s in this study). After cutting the CNT film into a designated pattern, we laminated the TFMS-doped CNT film onto the 2D/3D FACsPbI\textsubscript{3} perovskite film. The force during the lamination step should be carefully controlled to avoid the cracks in the perovskite film from which an electrical short circuit occurs (Figure S1). We noted that such phenomenon was not observed with methylammonium lead iodide (MAPbI\textsubscript{3}) and unique to 2D/3D FACsPbI\textsubscript{3}, probably due to difference in mechanical properties. After the transfer, we peeled off the nitrocellulose film and applied spiro-MeOTAD solution to improve the contact between the CNTs and the perovskite as well as enhancing the charge selectivity of the CNTs.

The four-probe measurement was used to measure the sheet resistance of the CNT films upon TFMS vapor doping. According to Figure 1b, the bare CNT film has a sheet resistance of approximately 153.9 \( \Omega \) sq\textsuperscript{-1}. We can observe that the sheet resistance gradually decreases with an increase in the TFMS vapor doping time. There is a particularly sharp decrease in the sheet resistance from 30 s to 40 s of the doping time, followed by almost saturation of the sheet resistance with the doping time longer than 40 s. The sheet resistance decreased from 153.9 to 150.5 (2.2\% decrease), 147.9 (2.7\% decrease), 139.1 (9.6\% decrease), 121.1 (21.3\%) and 120.5 \( \Omega \) sq\textsuperscript{-1} (21.7\%) for 10, 20, 30, 40 and 50 s of doping time, respectively. Although the decrease in the sheet resistance is not as great as the reported drop-cast doping method,\textsuperscript{20} this *ex-situ* vapor-
assisted doping method successfully induces doping of the top-laminated CNT films without damaging the layers underneath. Vis–NIR optical absorbance spectroscopy can also show the $p$-type doping effect of CNTs.\textsuperscript{26, 27} Figure 1c reveals that the optical transitions of the van Hove singularities ($E_{M11}^S$, $E_{S22}^S$, $E_{S11}^S$) were slightly quenched after TFMS vapor doping. This indicates that the inter-band optical absorptions were suppressed due to the shift of the Fermi level by the TFMS molecules. The uniformity of the doping was investigated by elemental distribution analysis in Figure S2. The uniform distribution of oxygen and sulfur originating from the TEMS confirms that the doping is microscopically uniform. Raman spectroscopy can also tell us the degree of doping effect. In our case, however, there was no apparent blue shift of the G nor G’ bands of the vapor-doped CNT films (Figure S3). This is probably ascribed to the mild effect of the vapor doping and there have been several cases in which the doping effect was not visible by Raman spectroscopy.\textsuperscript{27}
Figure 1. (a) Schematic illustrations showing ex-situ doping of carbon nanotube (CNT) by using vaporized trifluoromethyl sulfonamide (TFMS) followed by lamination transfer on perovskite thin films for device fabrication (b) Measured sheet resistance with different exposure time to TFMS vapor. (c) Vis-NIR spectra of bare CNT (black), 10 s vapor-doped CNT (red), 30 s vapor-doped CNT (blue), and 50 s vapor-doped CNT (green).

After investigating the vapor doping effect on CNT, we fabricated CNT-PSCs in a structural configuration of glass/indium-doped tin oxide (ITO)/SnO$_2$/FACsPbI$_3$/TFMS-doped CNT/spiro-MeOTAD (Figure 2a). In typical metal electrode-based devices, light transmitted through a perovskite layer is reflected back to the perovskite layer by the metal electrode to contribute to light harvesting efficiency and thus short-circuit current density ($J_{SC}$) of which contribution was estimated to be as high as 10%. In case of the devices based on a non-reflective CNT electrode, however, the transmitted light is subject to be lost by absorption of the CNT. Therefore, it is important to minimize the transmitted light through the perovskite layer to maximize the $J_{SC}$. The 2D/3D FACsPbI$_3$ has a lower bandgap of 1.48 eV with higher absorption coefficient than that of conventional MAPbI$_3$ (1.57 eV, Figure S4a). Furthermore, a superior charge carrier lifetime as evidenced by a longer photoluminescence (PL) lifetime enables utilization of a thicker 2D/3D FACsPbI$_3$ perovskite film (ca. 600 nm) than that of MAPbI$_3$ film (ca. 550 nm) that can further minimize the transmission loss (Figure S4b). Indeed, a bare CNT based 2D/3D FACsPbI$_3$ PSC showed a higher $J_{SC}$ (23.4 mA/cm$^2$) than that of a MAPbI$_3$ PSC (20.6 mA/cm$^2$, Table S1), mainly due to higher external quantum efficiency (EQE) at longer wavelength region (Figure S5). Nevertheless, the lower open-circuit voltage ($V_{OC}$) of 2D/3D FACsPbI$_3$ PSCs based on CNT may be related to energy level alignment between the 2D/3D FACsPbI$_3$ and CNT electrode.

Figure 2b shows current density-voltage ($J$-$V$) curves of PSCs based on CNTs without and with doping, and the measured photovoltaic parameters are summarized in Table 1. With 30 s
TFMS-doped CNT, the highest PCE of the device reached 17.56% ($J_{SC}$: 24.21 mA/cm$^2$, $V_{OC}$: 1.005V and FF: 0.722), whereas that of the device based on bare CNT showed 15.83% ($J_{SC}$: 23.35 mA/cm$^2$, $V_{OC}$: 0.985V and FF: 0.688). The obtained PCE with 30 s of TFMS doping is the highest among the reported PCEs from CNT top electrode-based PSCs (Figure S6). However, the doping time longer than 30 s decreased the device performance to 15.83% ($J_{SC}$: 22.67 mA/cm$^2$, $V_{OC}$: 0.978V and FF: 0.716) (Table 1). EQE spectra of the devices based on bare CNT and CNT doped by TFMS for 30 s are measured in Figure 2c. The device based on the CNT doped by TFMS shows slightly enhanced EQE between the wavelengths of 500 to 800 nm compared to the one based on bare CNT. Consequently, the integrated $J_{SC}$s were calculated to be 22.98 and 23.31 mA/cm$^2$ for the device based on bare CNT and CNT doped by TFMS for 30 s, respectively. The calculated $J_{SC}$s from the EQE spectra are well-correlated with the $J_{SC}$s measured from J-V curves with less than 4% discrepancy. Figure 2d-g show statistical distribution of photovoltaic parameters without and with TFMS doping (30 s). Upon doping of the CNT with TFMS, all the photovoltaic parameters improved; the $J_{SC}$ increased from 22.73±0.71 to 23.52±0.39 mA/cm$^2$ (Figure 2d, 3.4% increase), the $V_{OC}$ increased from 0.991±0.014 to 1.010±0.014 V (Figure 2e, 1.9% increase), and the FF increased from 0.679±0.019 to 0.706±0.016 (Figure 2f, 4.0% increase). As a result, the average PCE increased by 9.7% from 15.29±0.46 to 16.77±0.46% (Figure 2g).

Table 1. Photovoltaic performance of CNT top electrode-based 2D/3D FACsPbI$_3$ PSCs with different TFMS vapor doping time under 1 sun (AM 1.5 G, 100 mW/cm$^2$). The photovoltaic parameters are the best values with averages and error ranges in parentheses.

<table>
<thead>
<tr>
<th>Device ID</th>
<th>$J_{SC}$ (mA/cm$^2$) [average]</th>
<th>$V_{OC}$ (V) [average]</th>
<th>FF [average]</th>
<th>PCE (%) [average]</th>
</tr>
</thead>
<tbody>
<tr>
<td>Bare CNT</td>
<td>23.35 [22.73±0.71]</td>
<td>0.985 [0.991±0.014]</td>
<td>0.688 [0.679±0.019]</td>
<td>15.83 [15.29±0.46]</td>
</tr>
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Figure 2. (a) A schematic illustration showing structure of planar heterojunction perovskite solar cells based on TFMS doped CNT (b) Current density-voltage ($J$–$V$) curves of the highest efficiency perovskite solar cells incorporating bare CNT and CNT doped by TFMS (30 s and 50 s of doping time). (c) external quantum efficiency (EQE) spectra and corresponding integrated short-circuit current density ($J_{SC}$) of the device based on bare CNT and CNT doped by TFMS for 30 s (d) $J_{SC}$, (e) open circuit voltage ($V_{OC}$), (f) fill factor (FF) and (g) power conversion efficiency (PCE) of the perovskite solar cells incorporating bare CNT and CNT doped by TFMS.

To elucidate the reasons behind the improved photovoltaic parameters, we investigated the work functions of CNTs and resulting charge extraction properties by photoelectron yield spectroscopy (PYS) and photoluminescence (PL) analyses. PYS data of the TFMS-doped CNT films show that the Fermi level of CNT decreases with the increase in the TFMS vapor doping time (Figure 3a). As the doping time increased, the Fermi level of the CNT shifted from -4.75 eV (bare CNT) to -4.87 eV, -4.96 eV and -5.02 eV for 10 s, 30 s and 50 s doping time,
respectively. The higher work functions with the increased doping time can be attributed to the increase in hole concentration, owing to the larger amount of the TFMS adsorption, which is in agreement with the sheet resistance and Vis-NIR measurement in Figure 1. According to the energy diagram in Figure 3b, the higher work function caused by stronger TFMS doping is energetically favorable in reducing the potential loss in terms of hole extraction, which is probably the origin of the enhanced $V_{OC}$. The charge extraction dynamic was investigated by steady-state and time-resolved PL measurement in Figure 3c and d, respectively. The bare perovskite film in contact with a glass substrate showed the peak PL intensity of 4.58x10^6, which was significantly quenched to 1.07 x10^6 when in contact with bare CNT. With 30 s TFMS doping, the PL quenching became more pronounced where the peak PL intensity was 1.40 x10^6, implying better hole extraction capability of the CNT with TFMS doping. The better hole extraction ability after TFMS doping was further confirmed by time-resolved PL decay profiles (Figure 3d). The perovskite films in contact with CNT sheets showed the ultrafast decay followed by a slower decay in PL intensity. The ultrafast decay ($\tau_{HTM}$) can be attributed to the PL decay as a result of hole extraction by CNT, whereas the slower decay can be ascribed to the radiative recombination of remaining charges in bulk perovskite. The TFMS doped CNT showed the faster decay of the PL intensity ($\tau_{HTM}$=1.0 ns) with less proportion of remaining tail than that of the perovskite film with bare CNT ($\tau_{HTM}$=1.6 ns), indicating enhanced charge extraction capability with TFMS doping.

Despite the increase in the conductivity of CNT and energetically more favour outcome upon TFMS doping, the device performance decreased when the vapor doping time was longer than 30 s. We hypothesize that this is because of surplus TFMS molecules reacting with the spiro-MeOTAD solution. The spiro-MeOTAD solution contains a small amount of $t$-BP which is a
Lewis base with a nitrogen atom. TFMS can undergo an acid-base reaction with $t$-BP (1). This was evidenced by the downshifts of $H_a (\delta = 8.87 \text{ ppm})$ and $H_b (\delta = 8.16 \text{ ppm})$ in $t$-BP to $H_a (\delta = 8.49 \text{ ppm})$ and $H_b (\delta = 7.39 \text{ ppm})$ upon addition of 10 wt% TFMS, according to the $^1\text{H}$ NMR data in Figure 4. The change in the coupling constant indicates the change in chemical environment by the protonation (Figure S7). The protonated $t$-BP product precipitated in chlorobenzene as

Figure 3. (a) photoelectron yield spectroscopy (PYS) measurement data of bare CNT (black circles), 10 s vapor-doped CNT (red circles), 30 s vapor-doped CNT (blue circles), and 50 s vapor-doped CNT (green circles). The broken lines are the linear fit of the data. (b) Band alignment in the planar heterojunction perovskite solar cells without and with doping of the CNT (c) Steady-state photoluminescence (PL) spectra and (d) time-resolved PL decay profiles of bare perovskite film and perovskite films in contact with bare CNT and CNT doped by TFMS.
salts do not dissolve in non-polar solvents. The loss of t-BP from the spiro-MeOTAD solution leads to the decrease in all photovoltaic parameters for t-BP is absolutely necessary to inhibit the complexation and enhance the solubility of spiro-MeOTAD.\textsuperscript{31} This explains the decreases in $J_{SC}$, $V_{OC}$, and FF of the devices when longer than 30 s of TFMS doping was applied. However, the decreases in $V_{OC}$ and FF were not as big as the $J_{SC}$ drop. This is because while the loss of t-BP decreased all photovoltaic parameters, the decrease in the Fermi level and the increased conductivity of CNT countered the drops in $V_{OC}$ and FF. It should be mentioned that both spiro-MeOTAD and LiTFSI also contain nitrogen atoms, which may undergo protonation by TFMS as well.

$$4\text{-C(CH}_3\text{)}_3\text{C}_2\text{H}_4\text{N} + \text{CF}_3\text{SO}_3\text{H} \leftrightarrow [4\text{-C(CH}_3\text{)}_3\text{C}_5\text{H}_4\text{NH}]^+\text{[CF}_3\text{SO}_3\text{]}^- \quad (1)$$

This agrees with the Nyquist plot of 50 s TFMS-treated CNT showing higher counter electrode resistance ($R_{CT}$, first semi-circle) despite lower series resistance ($R_S$, starting point of the semi-circle) than either 30 s TFMS-treated CNT and bare CNT on a perovskite film (Figure S8). Therefore, it is of importance to control the amount of TFMS dopants. The optimal amount should be large enough to enhance the conductivity and to lower the Fermi level of the CNT, but small enough to avoid reaction with t-BP in the spiro-MeOTAD solution.
Figure 4. $^1$H NMR spectra were recorded in DMSO-$d_6$ solution at 500 MHz. α) $^1$H NMR of 4-tert-butylpyridine; β) $^1$H NMR of 4-tert-butylpyridine with addition of 10 wt% of trifluoromethanesulfonic acid.

Thermal- and light-stability test was conducted on bare CNT-PSCs, TFMS-doped CNT-PSCs, and the metal electrode-used PSCs as the reference. All three different types of devices were encapsulated under N$_2$ atmosphere and subjected to the open-circuit under constant illumination of 1 sun (the light was generated by halogen lamps) in air (60±5 °C, relative humidity of 50±10%). The results in Figure 5a show that the CNT top electrode-used PSCs are much more stable than the metal electrode-used CNT PSCs of which the degradation can be ascribed to migration of metal ions followed by reaction with perovskite.9, 18 With a metal electrode, the PCE of the device was degraded by 68.0% from the initial PCE, whereas the CNT based devices retained higher PCE after 244 h of illumination (83.3% of the initial PCE for the bare CNT-PSCs and 86.9% of the initial PCE for the TFMS-doped CNT-PSCs). It was interesting to observe that the TFMS-doped CNT-PSCs were slightly more stable than the bare CNT-PSCs. To investigate the reason behind this phenomenon, water contact angles were measured on CNT films with different vapor doping times (Figure S9). The results show that the water contact angle increases upon TFMS vapor doping. This denotes that the application of TFMS makes the CNT electrodes slightly more hydrophobic due possibly to the fluorinated end of TFMS sticking out from CNT.32 Although the difference is marginal, such an increase in hydrophobicity reduces the intrusion of moisture, which has a direct impact on the stability of the perovskite in a long-term operation. The contact angle decreased when the vapor doping time exceeded 30 s. We conjecture that this is because the charged acidic end of TFMS outweighing the effect of the fluorinated part of TFMS when an excess amount of TFMS is introduced.
Water vapor transmittance rate (WVTR) measurement was used to further investigate the water transmitting property. According to Figure 5b, the 30 s TFMS vapor-doped CNT film shows a slightly lower WVTR value than the bare CNT film. For the CNT film with the 50 s TFMS vapor doping, WVTR value is slightly higher than the bare CNT film, supporting our claim about the hydrophobicity of the CNT films. Although the WVTR value difference is small, the accumulated amount of transmitted water makes a substantial difference over a long period of time (inset of Figure 5b). This must have led to the difference in the stability between the TFMS-doped CNT-PSCs and the bare CNT-PSCs. We also suspect the charge extraction may have influenced the stability. As discussed above, the TFMS-doped CNTs have a better energy alignment with a lower $V_{OC}$ loss than the bare CNT-PSCs. Since not all of the uncollected charge gets recombined, delocalized charge within the perovskite system can accelerate the trapped charge-driven degradation.

In summary, we introduced an effective ex-situ vapor-assisted doping of a CNT top electrode for implementation of efficient and stable perovskite solar cells. The ex-situ TFMS vapor doping enabled enhancement of conductivity with a favorable alignment of work function without damaging underneath hole transporting and perovskite layers. Owing to enhanced conductivity with a favorable energy alignment, we achieved a PCE of 17.56% with enhanced stability from CNT electrode-laminated PSCs. It is worth mentioning that our method is simple, quick and readily applicable to large area module process. As the record-high PCE showcases, we successfully provided a solution to overcoming the efficiency limit in CNT-laminated PSCs. We believe this method can pave the pathway to higher performance and more stable CNT-incorporated PSCs for versatile applications such as flexible or stretchable devices.
Figure 5. (a) Evolution of power conversion efficiency (PCE) of the encapsulated devices under continuous one sun illumination generated by halogen lamps. The devices were under open-circuit condition (T=60±5 °C, relative humidity of 50±10%). (b) Water vapor transmittance rates (WVTRs) of bare CNT (black) film, 30 s vapor-doped CNT film (red), and 50 s vapor-doped CNT film (blue).

METHODS

Synthesis of CNT sheets and vapor-assisted doping. CNT films were synthesized by an aerosol (floating catalyst) CVD method based on ferrocene vapor decomposition in a CO atmosphere. The catalyst precursor was vaporized by passing through a cartridge filled with ferrocene powder. The flow containing ferrocene vapor was then introduced into the high-temperature zone of a ceramic tube reactor through a water-cooled probe and mixed with additional CO. To obtain stable growth of CNTs, a controlled amount of CO$_2$ was mixed with the CO carbon source. CNTs were directly collected downstream of the reactor by filtering the flow through a nitrocellulose or silver membrane filter (Millipore Corp., USA; HAWP, 0.45 μm pore diameter). The CNT films were vapor-doped by placing in the fume of TFMS for a designated time before transferring onto perovskite devices.
Materials Characterizations. Surface morphology and elemental distribution analyses of the CNT film was performed by scanning electron microscopy (SEM, Nova Nano 230). Steady-state photoluminescence (PL) spectrum was acquired by a Horiba Jobin Yvon system. A 640 nm monochromatic laser was used to excite the samples. Time-resolved PL decay profiles were recorded using a Picoharp 300 with time-correlated single-photon counting capabilities. The films were excited by a 640 nm pulse laser with a repetition frequency of 100 kHz provided by a picosecond laser diode head (PLD 800B, PicoQuant). The power of the excitation light was controlled to be ca. 1.4 nJ/cm$^2$ to minimize the carrier annihilation and non-geminate recombination. The Fermi levels were measured by Riken Keiki photoelectron yield spectroscopy in air (PYS-A) AC-2 and kelvin probe S spectroscopy in air (ESA). They were calibrated by Au before the measurement. WVTR measurement was carried out by Sumitomo Chemical Co. Ltd. Four probe was measured by a home-built system. Shimadzu UV-3150 was used for the UV-vis measurement. Solartron SI1287 Electrochemical Interface and Solartron 1255B Frequency Response Analyzer were used for the Impedance Measurement.

Device fabrication. Indium doped tin oxide (ITO) glass substrates were cleaned with successive sonication in detergent, deionized (DI) water, acetone and 2-propanol baths for 15 min respectively. The cleaned substrates underwent the UV-ozone for enhanced wettability. 30 mM SnCl$_2$·2H$_2$O (Aldrich, >99.995%) solution was prepared in ethanol (anhydrous, Decon Laboratories Inc.) as a precursor solution for deposition of a compact SnO$_2$ layer. The solution was filtered by 0.2 μm syringe filter, followed by spin-coating on the cleaned substrate at 3000 rpm for 30 s. The spin coated film was annealed at 150 °C for 30 min. After cooling down to room temperature, another cycle of the spin-coating process was repeated, which was followed by annealing at 150 °C for 5 min and 180 °C for 1 h. The SnO$_2$ coated ITO glass was treated with
UV-ozone before spin-coating of perovskite solution. The perovskite layer was prepared by the adduct approach. Typically, 163.0 mg of HC(NH$_2$)$_2$ (FAI), 453.4 mg of PbI$_2$ (TCI, 99.99%), 8.2 mg of PEAI, 5.0 mg of CsI (Alfa Aesar, 99.999%) and 97.4 mg of N-Methyl-2-pyrrolidone (NMP, Sigma-Aldrich, anhydrous, 99.5%) were added to 560 mg of N,N-Dimethylformamide (DMF, Sigma-Aldrich, anhydrous, 99.8%). Perovskite and hole transporting layer was formed in a glove box filled with dry air. The perovskite solution was spin-coated at 4000 rpm for 20 s to which 0.15 mL of diethyl ether (anhydrous, >99.0%, contains BHT as stabilizer, Sigma-Aldrich) was dropped after 10 s. The resulting transparent adduct film was heat-treated at 80 °C for 1 min followed by 150 °C for 20 min. The spiro-MeOTAD solution was prepared by dissolving 85.8 mg of spiro-MeOTAD (Lumtec) in 1 mL of chlorobenzene (anhydrous, 99.8%, Sigma-Aldrich) which was mixed with 33.8 μl of 4-tert-butylpyridine (96%, Aldrich) and 19.3 μl of Li-TFSI (99.95%, Aldrich, 520 mg/mL in acetonitrile) solution. The CNT film was laminated onto the perovskite film with rolling speed and pressure of approximately 0.5 cm/s and 0.01 N/cm$^2$, respectively. A supporting nitrocellulose film was removed after the lamination. The spiro-MeOTAD solution was applied onto the CNT layer to improve the contact between the CNTs and the perovskite. The spiro-MeOTAD solution was spin-coated on the perovskite layer at 3000 rpm for 20 s by dropping 17 μl of the solution during the spinning. For deposition of a metal electrode, an approximately 100 nm-thick silver layer was thermally evaporated at 0.5 Å/s on top of the spiro-MeOTAD layer.

Device characterizations. Current density-voltage ($J$–$V$) curves of the devices were recorded using Keithley 2401 source meter under simulated one sun illumination (AM 1.5G, 100 mW/cm$^2$) generated from Oriel Sol3A with class AAA solar simulator (Newport). The intensity of the light was calibrated by NREL-certified Si photodiode equipped with KG-5 filter. The $J$–$V$
curves were recorded at 0.1 V/s (between 1.2 V and -0.1 V with 65 data points and 0.2 s of delay time per point). The devices were covered with metal aperture with an area of 0.100 cm$^2$ to define the active area during the measurement. All the devices were measured without pre-conditioning (light-soaking and/or applied bias voltage. The external quantum efficiency (EQE) measurement system consisted of an MLS-1510 monochromator to scan the UV–vis spectrum. A source measurement unit was used to record the current at each specific wavelength.

ASSOCIATED CONTENT

The Supporting Information is available free of charge on the ACS Publications website.

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**Notes**

The authors declare no competing financial interest.

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